

An adaptive method for free energy computations

Kimiya Minoukadeh (CERMICS, ENPC)

Joint work with
Tony Lelièvre (CERMICS), Chris Chipot (University of Illinois)

Sampling the canonical measure

Consider a system of N particles with coordinates $q \in \mathbb{R}^{3N}$ interacting through the potential $V : \mathbb{R}^{3N} \rightarrow \mathbb{R}$.

In statistical physics we are often interested in sampling the **canonical measure**

$$\mu(dq) = Z^{-1} e^{-\beta V(q)} dq$$

where $\beta = 1/(k_B T)$.

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To sample this measure we use the **overdamped Langevin** dynamics

$$dX_t = -\nabla V(X_t) dt + \sqrt{2\beta^{-1}} dW_t$$

where $X_t \in \mathbb{R}^d$ is the system trajectory and W_t a standard Brownian motion.

Free Energy

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- Assume that the slow variable is $\xi(q) = z$, where $\xi : \mathbb{R}^{3N} \rightarrow \mathbb{R}$.
- ξ is called the **collective variable** or **reaction coordinate (RC)**.
- The image of the canonical measure in ξ has density

$$\psi^\xi(z) = \int_{\mathbb{R}^d} \psi(q) \delta_{\xi(q)-z}$$

- The **free energy** is then defined by $A(z) = -\beta^{-1} \ln \psi^\xi(z)$.

How can we efficiently compute the free energy?

Free Energy difference

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It can be shown that $A'(z)$ is the conditional expectation

$$A'(z) = \mathbb{E} \left[F^V(X) | \xi(X) = z \right]$$

$$\text{where } F^V = \frac{\nabla V \cdot \nabla \xi}{|\nabla \xi|^2} - \beta^{-1} \operatorname{div} \left(\frac{\nabla \xi}{|\nabla \xi|^2} \right).$$

Adaptive Biasing Force methods

The idea of the **Adaptive Biasing Force** (ABF) method is to bias the standard dynamics:

$$\begin{cases} dX_t = -\nabla(V - A_t \circ \xi)(X_t)dt + \sqrt{2\beta^{-1}}dW_t \\ A'_t(z) = \mathbb{E} \left[F^V(X_t) | \xi(X_t) = z \right] \end{cases}$$

If the process X_t has law $\psi_t dq$ then the marginal law in ξ satisfies

$$\partial_t \psi_t^\xi = \beta^{-1} \partial_{zz} \psi_t^\xi$$

Furthermore, if $A'_t = A'$, then $\psi_\infty = Z^{-1} e^{-\beta(V - A \circ \xi)}$ and the marginal measure in ξ is uniform.

Adaptive Biasing Force method

The aim of the ABF method is therefore twofold:

- 1 to serve as an **adaptive importance sampling** method.
- 2 to compute **free energy differences** efficiently.

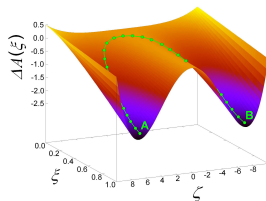
But in practice, how is A'_t computed?

Single-replica ABF

- 1 Original ABF: trajectorial averaging (Darve *et al.*, 2001)

$$A'_t(z) = \frac{\int_0^t F^V(X_s) \delta_{\xi(X_s)-z} ds}{\int_0^t \delta_{\xi(X_s)-z} ds}$$

- Single-replica simulations can lead to **slow convergence**
- We may be overlooking **other slow degrees of freedom**



Multiple-replica ABF

Use of multiple replicas can speed up convergence (Lelièvre et al., 2007)

- ② MR-ABF: averaging over trajectories and $R > 1$ replicas

$$A'_t(z) = \frac{\sum_{i=0}^{R-1} \int_0^t F^V(X_s^i) \delta_{\xi(X_s^i)-z} ds}{\sum_{i=0}^{R-1} \int_0^t \delta_{\xi(X_s^i)-z} ds}$$

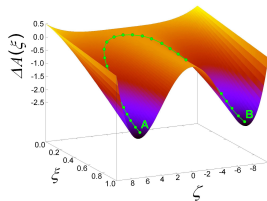
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- Replicas are likely to explore different valleys
- Method is easily parallelized...



Selection: improving sampling in ξ

Advantage of MR-ABF: we can further accelerate the diffusion in ξ :

$$\partial_t \psi_t^\xi = \beta^{-1} \partial_{zz} \psi_t^\xi + S(t, z) \psi_t^\xi$$

Choosing $S(t, z) = c \frac{\partial_{zz} \psi_t^\xi(z)}{\psi_t^\xi(z)}$ gives rise to

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This can be implemented by assigning each replica i with weight

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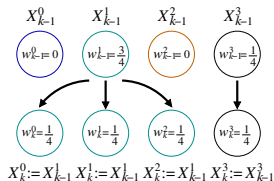
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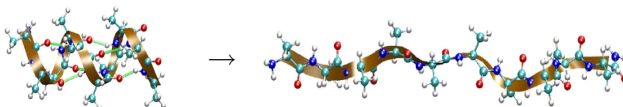
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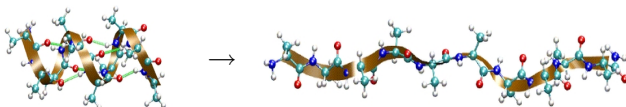
Application: deca-alanine

These two methods are compared for the **unfolding of the deca-alanine peptide**



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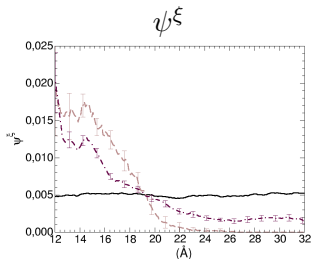
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The reaction coordinate ξ is chosen as the end-to-end distance of the peptide chain.

- From α -helix to extended states $\xi : 12 - 32 \text{ \AA}$
- Compact states $\xi : 4 - 16 \text{ \AA}$

Extended conformations

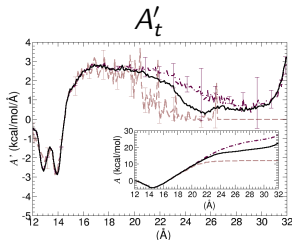


ξ : 12 – 32 Å



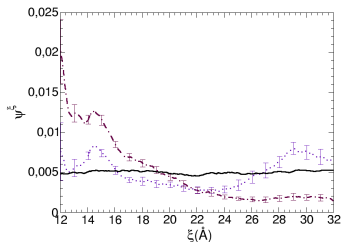
Results from simulations after 0.25 ns:

- **Single-replica** simulations rarely stretch beyond $\xi = 22$ Å
- **Multiple-replica** simulations explore the whole RC space and mean force approximations nearly converged



- Original ABF (1 replica)
- . - . MR-ABF (16 replicas)
- Reference

Results in range ξ : 12 – 32 Å

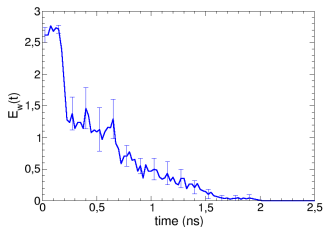


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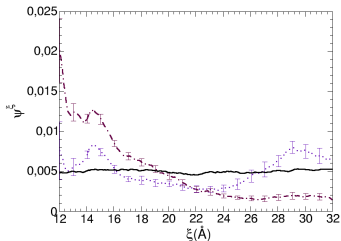


Results from 16-replica simulations after 0.25 ns

- · — · MR-ABF
- · · · MR-ABF with selection
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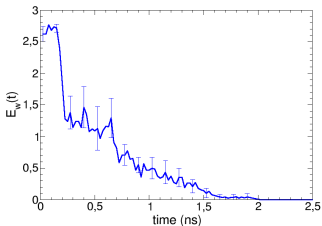


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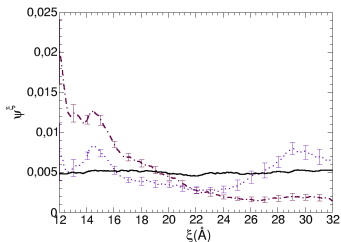
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Relative entropy of weights:

$$E_w(t) = \sum_{i=0}^{R-1} w_t^i \log(w_t^i) - \log(1/R)$$

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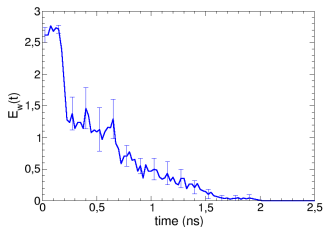


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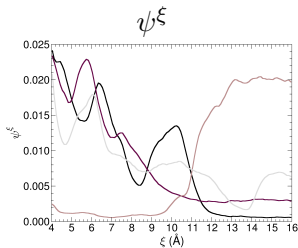


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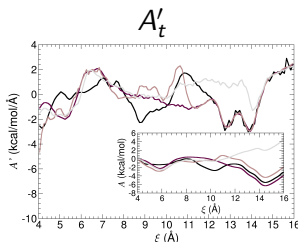
$$E_w(t) = \sum_{i=0}^{R-1} w_t^i \log(w_t^i) - \log(1/R)$$

Stop selection when $E_w(t) < \varepsilon \log(R)$

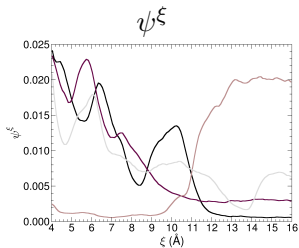
Compact conformations: 1-replica ABF



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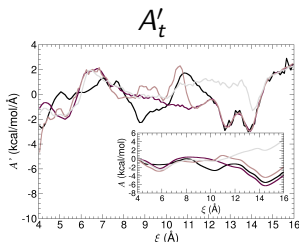


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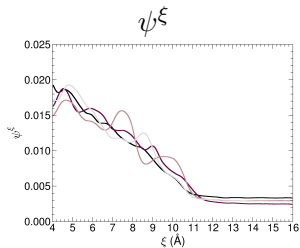
Results from four independent **single-replica** simulations after **100 ns**:

- Mean force estimations are inconsistent
- One simulation shows global minimum at $\xi = 6 \text{ \AA}$



Why? Metastabilities at fixed ξ .

Compact conformations: 32-replica ABF

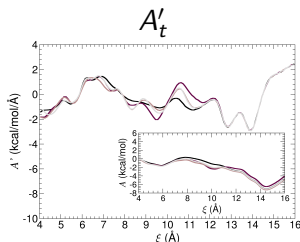


$\xi : 4 - 16 \text{ \AA}$



Results from four independent **32-replica** simulations after **100 ns**:

- Qualitatively consistent mean force approximations
- The α -helix conformation recovered as the global minimum



Conclusion

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Multiple replicas speed up convergence of ABF

- **well chosen RC**: selection can accelerate diffusion in ξ .
- **suboptimal RC**: multiple replicas help to explore parallel valleys

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Current work

Mathematical proof of the convergence of MR-ABF in the case of multiple valleys.

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Reference: K. Minoukadeh, C. Chipot and T. Lelièvre, *Parallel algorithms for free energy calculations: a multiple replicas adaptive biasing force approach*. In preparation